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**Citation:** Gulistan, A., Ghosh, S., Ramachandran, S. & Rahman, B. M. (2017). Efficient strategy to increase higher order inter-modal stability of a step index multimode fiber. *Optics Express*, 25(24), pp. 29714-29723. doi: 10.1364/OE.25.029714

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# Efficient strategy to increase higher order inter-modal stability of a step index multimode fiber

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**Abstract:** We demonstrate a novel approach to enhance the mode stability through increased effective index difference ( $\Delta n_{eff}$ ) between the higher-order modes ( $LP_{18}$ ,  $LP_{09}$  and  $LP_{19}$ ) of a multimode fiber. Fibers with large diameters have bigger effective mode areas ( $A_{eff}$ ) and can be useful for high power lasers and amplifiers. However, a large mode area (LMA) results in an increased number of modes that can be more susceptible to mode coupling. The modal effective index difference ( $\Delta n_{eff}$ ) strongly correlates with mode stability and this increases as the modal order ( $m$ ) increases. We report here that the mode spacing between the higher order modes can be further enhanced by introducing doped concentric rings in the core. In our work, we have shown a more than 35% increase in the mode spacing between the higher order modes by optimizing the doping profile of a LMA fiber. The proposed design technique is also scalable and can be applied to improve the mode spacing between different higher order modes and their neighboring antisymmetric modes, as necessary.

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**OCIS codes:** (060.2400) Fiber properties; (060.2310) Fiber optics; (230.2285) Fiber devices and optical amplifiers; (060.2280) Fiber design and fabrication.

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## 1. Introduction

High power lasers and amplifiers are necessary for many applications in the field of science and technology. Recent research shows remarkable increase in the output power of fiber lasers [1, 2]. Compared to solid state lasers, fiber lasers are more compact, lightweight and flexible. Various approaches have been considered to increase the power of fiber based lasers and amplifiers. However, an increase in power brings the need to mitigate for non-linearities in the fiber such as Brillouin and Raman scattering, four-wave mixing (FWM) and self-phase modulation [3, 4]. Mitigating these non-linear effect require large mode area (LMA) fibers in which optical mode is less intense. Therefore, much of the recent research focuses on LMA fibers for the high-power fiber lasers [5]. Multimode fibers (MMF) offer large mode areas and high beam quality [6] but the existence of many modes can result in random mixing between the desired mode and unwanted modes. However, recent research shows that higher order modes (HOM) have the ability to provide for more stable operation because the signal stability increases with an increase in the modal order [5], [7, 8]. A number of approaches have been presented in recent years to use HOM for high power fiber lasers and amplifiers. One of the approaches involves the coupling of light from the fundamental mode to a single desired HOM using fiber Bragg gratings and then to propagate light in the higher order mode amplifying fiber [9]. These techniques provide considerably larger modal areas as well as more stable operation as described above. Moreover, for a given effective area, HOMs are also less sensitive to area reduction than the fundamental mode when bending [10, 11]. At the same time, compared to the fundamental mode, higher order modes are less prone to mode coupling as with the increase in the modal order ( $m$ ) the difference in effective index ( $\Delta n_{\text{eff}}$ ) between a given higher mode ( $LP_{0m}$ ) and its nearest antisymmetric mode ( $LP_{1m}$ ) also increases [12].

In this paper, we have proposed a novel approach to increase the  $\Delta n_{\text{eff}}$  that would result in a resilience to mode-mixing and more stable signal propagation with the advantage of significantly larger effective mode area ( $A_{\text{eff}}$ ) by using HOMs. Annular rings with doping of increased or reduced refractive index are used at particular radial locations inside the core as shown in Fig. 1, such that the effective refractive index ( $n_{\text{eff}}$ ) of a desired mode is increased or reduced [13]. To demonstrate our concept, we have considered a MMF of numerical aperture (NA=0.22) with core radius ( $R_{\text{core}}=25 \mu\text{m}$ ) and refractive index ( $n_{\text{core}}=1.457$ ). Similarly, cladding radius and refractive index are taken as ( $R_{\text{clad}}=62.5 \mu\text{m}$ ) and ( $n_{\text{clad}}=1.4403$ ), respectively. For this study the central operating wavelength of ( $\lambda=1.05 \mu\text{m}$ ) is chosen.

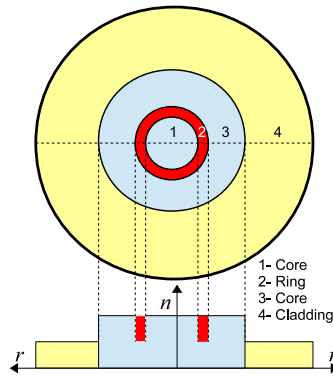


Fig. 1. Ring doping schematic of a MMF with the change in the refractive index along r-axis.

## 2. Theory

### 2.1. Modal solutions

Use of the higher order mode (HOM) of a multimode fiber (MMF) shows several advantages which includes mode area scaling to control laser high power and dispersion management for ultra-short pulses. Modal instabilities in HOMs are a common problem arising due to external perturbations, such as bending and fabrication imperfections. According to the symmetry rule, bend perturbation is odd in nature and direct coupling arises between even and odd order modes. Among the many modes guided by a MMF, sometimes a desired dominant mode ( $LP_{0m}$ ) may transfer its energy to its nearest anti-symmetric ( $LP_{1(m+1)}$  and  $LP_{1(m-1)}$ ) modes on both sides [12]. A low effective index difference ( $\Delta n_{eff}$ ) between adjacent modes enhances the modal energy transfer which results actual modal energy loss along with the interference effects. This could result in inter-mode mixing in a MMF. Although a lower order (lower value of  $m$ ) mode may be easier to excite, a higher value of  $m$  gives larger modal separation ( $\Delta n_{eff}$ ) values, so a compromise is needed. However, for a given value of  $m$ , if the modal separation to the nearest anti-symmetric modes can be increased, this would be a more preferable design. Here, a novel MMF design with several concentric material strips at strategically located positions is proposed to increase the  $\Delta n_{eff} = (n_{eff}^{0m} - n_{eff}^{1m})$  in desired modes, such as,  $LP_{18}$ ,  $LP_{09}$  and  $LP_{19}$  as a high  $\Delta n_{eff}$  would restrict unwanted modal energy transfer due to external perturbations. Instead of using a perturbation approach, we have used a rigorous full vectorial  $\mathbf{H}$ -field based finite element method (FEM) [14, 15] to find the modal solutions of our proposed MMF design. The FEM is one of the most numerically efficient and accurate approaches to obtain the modal solutions of an optical waveguide, to calculate the propagation constant ( $\beta$ ), effective index ( $n_{eff} = \beta/k_o$ , where wavenumber  $k_o$  is given by  $2\pi/\lambda$ ) and  $A_{eff}$ .

In order to increase the modal solution accuracy, the available two-fold symmetry of the fiber is exploited and only one-quarter of the structure is simulated [16, 17]. This allows more dense mesh distribution in the quarter structure of the fiber instead of distributing the same mesh over the whole structure. The polar mesh [18] discretization is also used, which can accommodate the discretized elements more efficiently at the circular boundaries, which can provide more accurate results compared to the mesh distribution based on the Cartesian coordinate system [19]. It is well known that the simulation accuracy of the FEM is highly dependent on the number of discretized elements used. Variation of the effective index ( $n_{eff}$ ) with the number of mesh elements ( $N$ ) is shown in Fig. 2 by a solid black line for the higher order  $LP_{09}$  mode. It can be observed that initially as mesh density increases the  $n_{eff}$  also increases rapidly and then asymptotically settles to a constant value. It should be noted that accuracy is up to the  $3^{rd}$  decimal place when the

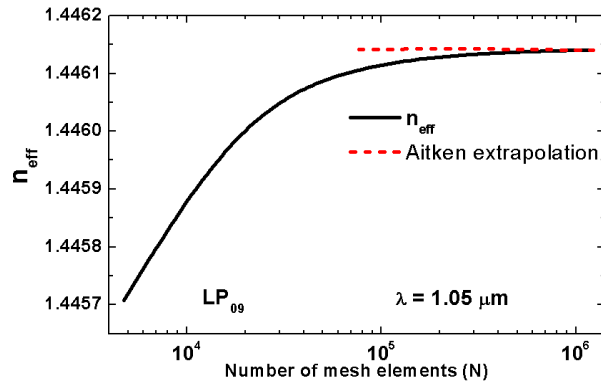


Fig. 2. Variation of  $n_{eff}$  of the  $LP_{09}$  mode with the mesh number ( $N$ ) and convergence realized with the Aitken extrapolation technique.

number of elements used is,  $N=7 \times 10^4$  and it increases to the 5<sup>th</sup> decimal place when  $N=5 \times 10^5$ . Here  $N$ , is the number of triangular elements used to represent a quarter of the MMF. A powerful Aitken's extrapolation technique is used to test the convergence of the modal solutions [20, 21]. Three successive  $n_{eff}$  values for corresponding mesh divisions with a geometric ratio are used in the Eq. (1).

$$n_{eff}^{\infty} = n_{eff(r+1)} - \frac{[n_{eff(r+1)} - n_{eff(r)}]^2}{n_{eff(r+1)} - 2n_{eff(r)} + n_{eff(r-1)}} \quad (1)$$

Using Eq.1 the extrapolated values of  $n_{eff}^{\infty}$  are shown in the Fig.2 by the red-dashed line. Aitken's values are calculated for the mesh values  $N=7.66 \times 10^4$ ,  $3.07 \times 10^5$ , and  $1.23 \times 10^6$  increased in fixed geometric ratio yielding the  $n_{eff}$  values 1.4460355, 1.4461157, and 1.4461358, respectively. It should be noted that in each solution the mesh density is two times (no. of the elements is four times) that of the previous and thus the geometrical mesh ratio is kept constant. From these values a more accurate extrapolated value of 1.4461398 is obtained. Similarly, Fig. 2 clearly shows the convergence of the extrapolated results and raw FEM results. As the trend of  $n_{eff}$  with increasing  $N$  for different modes is similar, so the accuracy of  $\Delta n_{eff}$  between two modes with increasing  $N$  is also greater.

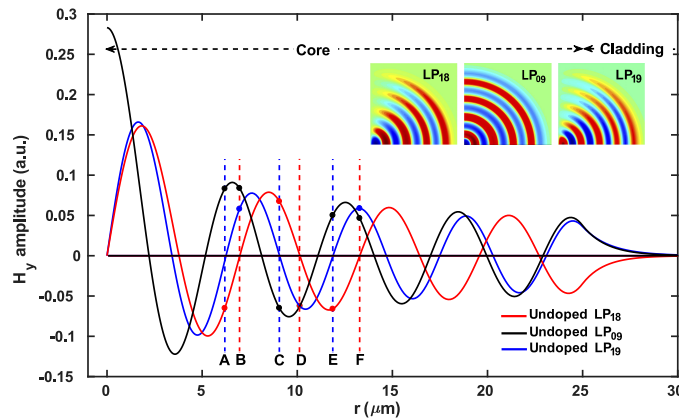


Fig. 3. Variations of  $H_y$  fields of the  $LP_{18}$ ,  $LP_{09}$ , and  $LP_{19}$  modes along the  $r$ -axis of MMF, contour field profiles in inset and the key points of interest are also shown.

## 2.2. Mode stability

As previously discussed, higher order modes (HOM) provide the unique benefit of increased stability due to a greater mode spacing ( $\Delta n_{eff}$ ) compared with that of the fundamental modes ( $LP_{01}$ ) [22]. In this paper, we propose a unique approach that can reduce mode mixing between the higher order  $LP_{18}$ ,  $LP_{09}$  and  $LP_{19}$  modes, that could alternatively only be possible by considering a much higher modal order. For our design, we used  $LP_{09}$  as the central propagation mode and modified the refractive index profile of the fiber such that the mode stability between symmetric ( $LP_{09}$ ) and the two nearest antisymmetric ( $LP_{18}$  and  $LP_{19}$ ) modes is increased. A measure of the modal stability can be defined using the effective index differences as,  $S_1 = \Delta n_{eff}(LP_{18} - LP_{09})$  and  $S_2 = \Delta n_{eff}(LP_{09} - LP_{19})$ . However, if required any other HOM of interest can also be selected and using the same concept its mode stability can be enhanced. For this structure, the effective indices of the  $LP_{18}$ ,  $LP_{09}$ , and  $LP_{19}$  are calculated as 1.447317, 1.4461104, and 1.4448554, respectively, yielding  $S_1 = 0.0012066$  and  $S_2 = 0.001255$ . The variation of the dominant  $H_y$  field profile of the  $LP_{18}$ ,  $LP_{09}$  and  $LP_{19}$  modes is shown in Fig. 3 by red, black and blue lines, respectively.  $H_y$  field contours of the  $LP_{18}$ ,  $LP_{09}$  and  $LP_{19}$  modes are also shown Fig. 3 as insets. It can be observed that the field profile of  $LP_{09}$  shown by a black line, has the highest magnitude at the center ( $r=0$ ) of the fiber and eight zero values along the radial direction. The antisymmetric  $LP_{18}$  and  $LP_{19}$  modes shown with red and blue lines, respectively, have zero field values at the center ( $r = 0$ ) of the fiber. The dominant  $H_y$  fields of the  $LP_{18}$  and  $LP_{19}$  modes have eight and nine zero field values along the radial direction, respectively. Table 1 shows all the radial locations where these modes have zero field values. For example,  $LP_{18}$  has its first three zero field locations at  $r=0$ , 3.78 and  $6.966 \mu\text{m}$ . Using these zero crossings we have identified multiple locations that are suitable for doping such that the mode stability ( $S_1$  and  $S_2$ ) can be increased.

Table 1. Zero crossing locations of field profiles of the  $LP_{18}$ ,  $LP_{09}$  and  $LP_{19}$  modes along r-axis ( $\mu\text{m}$ ).

Mode	Location of zero crossings along r-axis ( $\mu\text{m}$ )							
<b><math>LP_{18}</math></b>	0	3.78	6.966	10.138	13.279	16.315	19.557	22.684
<b><math>LP_{09}</math></b>	2.16	5.16	8.125	11.078	13.99	16.89	19.9	22.79
<b><math>LP_{19}</math></b>	0	3.36	6.193	9.058	11.867	14.577	17.472	20.252

Some specific points, A, B, C, D, E and F are selected as shown in Fig. 3, where either  $LP_{18}$  or  $LP_{19}$  has zero crossing. The reason for choosing these points is that we want to have less of an effect on one of the mode and have more of an effect on the other two though using strips of different doping. The modal field values at these points are also given in Table 2. For example at

Table 2. Field values of  $LP_{18}$ ,  $LP_{09}$  and  $LP_{19}$  at A, B, C, D, E and F points.

Mode	A=6.193	B=6.966	C=9.058	D=10.138	E=11.867	F=13.279
<b><math>LP_{18}</math></b>	0.06507	0	0.06741	0	0.6606	0
<b><math>LP_{09}</math></b>	0.08321	0.08328	0.06493	0.06218	0.05022	0.04658
<b><math>LP_{19}</math></b>	0	0.058	0	0.06267	0	0.05879

point A, the  $H_y$  field value of  $LP_{19}$  is zero, whereas, the field values of  $LP_{18}$  and  $LP_{09}$  modes are 0.06507 and 0.08321, respectively. As a result any change of refractive index doping at point A will have an almost negligible effect on the  $LP_{19}$  mode and comparatively more of an effect on the  $LP_{09}$  than the  $LP_{18}$  mode. However, at point F, where the  $H_y$  field value of  $LP_{18}$  mode is zero and these values for  $LP_{09}$  and  $LP_{19}$  modes are 0.04658 and 0.05879, respectively. Hence, at point F, doping will have no effect on the  $LP_{18}$  but will have more of an effect on  $LP_{19}$  than the



$LP_{09}$  mode. This will result in an increase in effective index both of these modes while keeping the effective index of the  $LP_{18}$  mode unchanged, and a suitable selection can result in an increase in the modal stability. In order to increase the  $\Delta n_{eff}$  between these modes a circular strip of  $0.3 \mu m$  width is considered that can have an increased or reduced refractive index by  $\Delta n$ . Here, we have taken  $\Delta n = 0.0167$ , which is also the difference between the core and cladding refractive indices. However, different values of  $\Delta n$  in these strips can be chosen according to the required level of stability between modes.

### 3. Numerical results

Table 3 shows the effect of doping a single individual layer at the above mentioned six positions along the r-axis. The second column of Table 3 shows the original  $\Delta n_{eff}$  between the modes without any doping. It can be noted that the  $\Delta n_{eff}$  between  $LP_{09}$  and  $LP_{19}$  is slightly higher than the  $\Delta n_{eff}$  between  $LP_{18}$  and  $LP_{09}$ . The values of  $+\Delta n$  or  $-\Delta n$  on particular points are chosen such that the effect on the central mode, in our case  $LP_{09}$ , is negligible or can be compensated with another doping layer where  $\Delta n$  is chosen with an opposite sign to the first point. For example, in layer A (at  $r = 6.193 \mu m$ ),  $+\Delta n$  is chosen to increase the stability,  $S_1$  between  $LP_{09}$  and  $LP_{19}$ , but unfortunately this reduces the stability,  $S_2$  between  $LP_{18}$  and  $LP_{09}$ . On the other hand a reduction of refractive index in layer B increases modal separation  $S_2$ , but reduces that of  $S_1$ . However, an increase of refractive index in layer E and reduction in layer F enhances both the  $S_1$  and  $S_2$  modal stabilities. The same doping approach is considered at all six positions with either  $+\Delta n$  or  $-\Delta n$  as shown in Table 3. Here, two approaches can be considered;

1. Using single layer doping to increase the stability between modes.
2. Using the combination of two or more layers to increase the stability.

Table 3. Individual strip doping effect on  $\Delta n_{eff}$  at points A, B, C, D, E, and F.

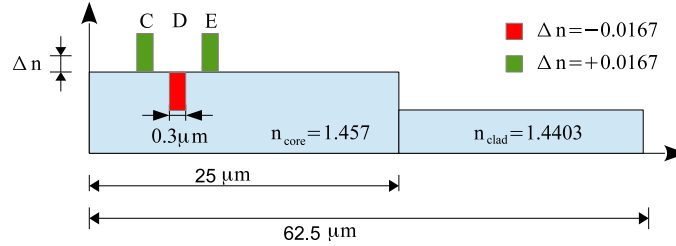
$\Delta n_{eff}$	Without doping	A	B	C	D	E	F
		$+\Delta n$	$-\Delta n$	$+\Delta n$	$-\Delta n$	$+\Delta n$	$-\Delta n$
$S_1 = LP_{18} - LP_{09}$	0.0012066	0.0010515	0.0015574	0.0012192	0.0014932	0.0013553	0.0014033
$S_2 = LP_{09} - LP_{19}$	0.001255	0.0015544	0.0010634	0.0015037	0.0012711	0.0014385	0.0014217

In Table 4 the above described approaches are shown with the percentage increase in  $\Delta n_{eff}$  ( $S_1$  and  $S_2$ ). The percentage increase is calculated with respect to the original  $\Delta n_{eff}$  between the modes as shown in column two. Here, three different options are suggested depending upon the required increase in the  $\Delta n_{eff}$ . It can be seen that with a single layer of  $-\Delta n$  doping at point F,  $S_1$  and  $S_2$  are increased by 16% and 14%, respectively. However, using two layers (at E & F points) the stabilities  $S_1$  and  $S_2$  between modes can be increased by 20% and 23%, respectively. It should be noted that the  $\Delta n$  doping at points E and F are taken as positive and negative, respectively. For further enhancement, three layers can be doped simultaneously at points C, D, and E which results in an increase of 35% and 38% for  $\Delta n_{eff}(LP_{18} - LP_{09})$  and  $\Delta n_{eff}(LP_{09} - LP_{19})$ , respectively. Hence, our proposed design results in increased modal spacing between the higher order modes  $LP_{18}$ ,  $LP_{09}$  and  $LP_{19}$ , thus providing for more stable and mode-mixing resistant operation. Figure 4 shows the refractive index profile for three layer doping at C, D and E points. Here,  $\Delta n$  at points C and E are taken as positive such that it increases the local refractive index from 1.457 to 1.4737. Whereas, at point D,  $\Delta n$  is taken as negative resulting in the local index being equal to that of the cladding. The combination is chosen such that  $\Delta n_{eff}(LP_{18} - LP_{09})$  and  $\Delta n_{eff}(LP_{09} - LP_{19})$  have an almost equal increase. Here, all three annular strips are centered at points C, D and E and have an equal width of  $0.3 \mu m$ . The  $LP_{09}$  mode is considered as a central

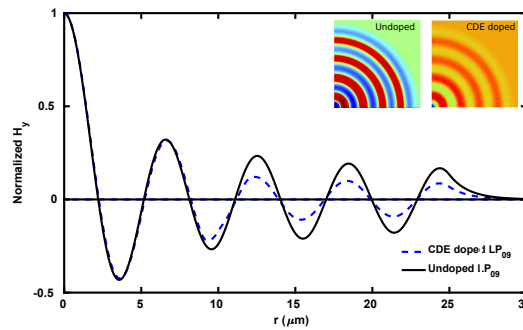


Table 4. Percentage increase in the  $\Delta n_{eff}$  using individual and combination approach.

$\Delta n_{eff}$	Without doping	F only		E & F		C, D & E	
		$\Delta n_{eff}$	% Increase	$\Delta n_{eff}$	% Increase	$\Delta n_{eff}$	% Increase
$S_1 = LP_{18} - LP_{09}$	0.0012066	0.0014033	16	0.0014551	20	0.0016254	35
$S_2 = LP_{09} - LP_{19}$	0.001255	0.0014217	14	0.0015318	23	0.0017357	38

Fig. 4. Refractive index profile of the modified MMF along  $r$ -axis with  $\pm \Delta n$  at C, D and E points.

propagation mode in our design. Hence next, the effect of three layers of doping at points C, D, and E points on the field profile of  $LP_{09}$  is studied. Figure 5 shows the  $H_y$  field profile of the  $LP_{09}$  mode before and after doping. The black line shows the undoped field profile whereas the dotted blue line represents the field profile of  $LP_{09}$  after doping these three layers at C, D, and E. It can be observed that until the appearance of the doped strips the field profile was almost unchanged, however beyond these strips, the field value is reduced compared to that of the original undoped fiber.

Fig. 5. Variations in the  $H_y$  fields of  $LP_{09}$  modes along the  $r$ -axis of the undoped fiber and the fiber with C, D, and E layers doped. The contour field profiles are also shown inset.

The  $\Delta n = 0.0167$  value used for the circular strips C, D, and E is equivalent to the refractive index difference between core and cladding of the fiber but of necessary other  $\Delta n$  values or even unequal values can be used for a particular design. To observe the effect of increased or reduced value of  $\Delta n$ , we have halved the refractive index difference as,  $\pm \Delta n = 0.00835$  and found that the modal stabilities  $S_1$  and  $S_2$  reduces to 18% and 19%, respectively. However, when refractive index difference is increased to double; as  $\pm \Delta n = 0.0334$ , the modal stabilities  $S_2$  increases to 74% but  $S_1$  increased only slightly to 41%, as field profiles were distorted significantly. However, it can be stated that a significant enhancement in the modal stability can be achieved by this approach.

Although we have focused on the enhancement of mode spacing between  $LP_{09}$  and its neighboring antisymmetric  $LP_{18}$  and  $LP_{19}$  modes, however, our approach is scalable and can be

applied for any higher order modes. As an example, for the enhancement of mode spacing of  $LP_{08}$  mode and its neighboring antisymmetric  $LP_{17}$  and  $LP_{18}$  modes, we have carried out additional simulations after introducing the similar annular rings. Similar six zero crossing points,  $A'$  to  $F'$  for  $LP_{17}$  or  $LP_{18}$  are identified as  $A'=6.928$ ,  $B'=7.856$ ,  $C'=10.138$ ,  $D'=11.518$ ,  $E'=13.279$ ,  $F'=15.098$  (in  $\mu\text{m}$ ). Numerical simulations were carried out with strips width of  $0.3 \mu\text{m}$  and  $\pm\Delta n=0.0167$  introduced at these points. Table 5 shows the increase in  $\Delta n_{eff}(LP_{17} - LP_{08})$  and  $\Delta n_{eff}(LP_{08} - LP_{18})$  with single layer ( $F'$ ) or multiple layers ( $E'$ ,  $F'$  and ( $C'$ ,  $D'$ ,  $E'$ ) with perturbed annular strips. It can be observed that the stability is increased to 46% for the  $LP_{08}$  mode when three annular layers at ( $C'$ ,  $D'$ ,  $E'$ ) points are used. This confirms that the concept presented here can be applied to any higher order modes, as necessary.

Table 5. Percentage increase in the  $\Delta n_{eff}$  of  $LP_{08}$  mode and its neighboring antisymmetric modes using individual and combination of two or three strips doping.

$\Delta n_{eff}$	Without doping	$F'$ only		$E'$ & $F'$		$C'$ , $D'$ & $E'$	
		$\Delta n_{eff}$	% Increase	$\Delta n_{eff}$	% Increase	$\Delta n_{eff}$	% Increase
$LP_{17}-LP_{08}$	0.0010746	0.0012218	14	0.0013424	25	0.0015759	47
$LP_{08}-LP_{18}$	0.0011146	0.0013136	18	0.0015187	36	0.0016256	46

### 3.1. Fabrication tolerance of strips width

Here, we demonstrate the effect of possible variations in doping that can occur during the fabrication process. As the combined doping of the three layers at points C, D, and E could be more sensitive to fabrication tolerances than the two or single layer doping we will consider the three layers (C, D and E) case for further investigation. The effects of a change in the layer width ( $w$ ) from  $0.3 \mu\text{m}$  to a higher or lower value are shown in Fig. 6. As discussed earlier, with  $w=0.3 \mu\text{m}$  the stabilities  $S_1$  and  $S_2$  between the modes are 35% and 38%, respectively. With an increase in the width from  $w=0.3 \mu\text{m}$  to  $w=0.4 \mu\text{m}$ , the stabilities  $S_1$  and  $S_2$  further increase to 38% and 51%, respectively. This is because the area of the doped layer is increased when the width is changed from  $0.3 \mu\text{m}$  to  $0.4 \mu\text{m}$  thus increasing its overall effect. However, as the width decreases to  $w=0.20 \mu\text{m}$ , the stability improvement reduces but still it remains above 25%. Hence, in our proposed design the modal stability improvement will remain at least 25% larger for a width change of  $\pm 0.1 \mu\text{m}$ .

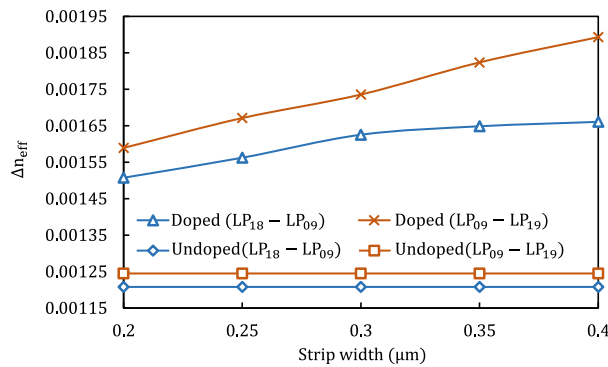


Fig. 6. Effect on  $\Delta n_{eff}$  of a change in width of doped layers at points C, D, and E.

### 3.2. Fabrication tolerance of strips center location

In our proposed design three layers are centered at points  $C=9.058 \mu\text{m}$ ,  $D=10.138 \mu\text{m}$  and  $E=11.867 \mu\text{m}$  and during fiber drawing these positions may change. Figure 7 shows the effect of a change in the location of the doped layer's center points. It can be seen that the stability is highly dependent on the specified locations in Fig. 3. When there is no shift in the location of the strip's centers, the stability improvement is larger than 35%. With a tolerance of  $\pm 0.05 \mu\text{m}$  the stability improvement still remains above 27% but as the shift is increased to  $0.1 \mu\text{m}$  the  $S_1$  improvement drops to 20% while  $S_2$  improvement is increased to 51%. Moreover, when the annular strips shift is  $-0.1 \mu\text{m}$ , the modal stability  $S_1$  improvement increases to 46% but that of  $S_2$  decreases to 27%. For comparison the spacing between the modes before introduction of the C, D and E layers is also shown by two horizontal lines.

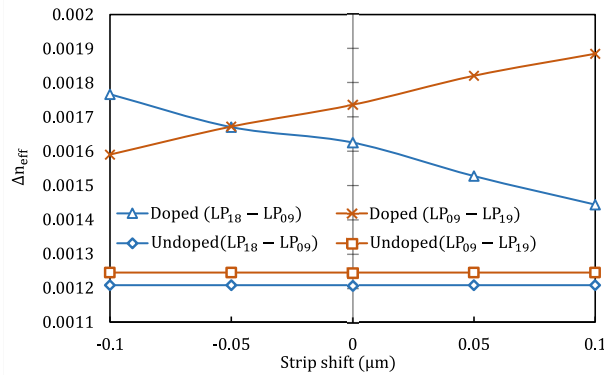


Fig. 7. Effect on the  $\Delta n_{eff}$  of a variation in the position of C, D, and E layers from center location.

### 3.3. Fabrication tolerance of wavelength change

The above analysis is carried out at the center wavelength of  $\lambda=1.05 \mu\text{m}$ . However, when the operating wavelength changes from this value the stability between modes can also change. To observe the impact of wavelength change on the stability between the modes, we varied the operating wavelengths and this is shown in Fig. 8. It can be observed that the stability between the  $LP_{18}$ ,  $LP_{09}$  and  $LP_{19}$  modes increases with an increase in the wavelength. It should be noted that the refractive index of a material is also dependent on the wavelength. In our analyses we have used the core and clad refractive index values as  $n_{core}=1.457$  and  $n_{clad}=1.4403$  respectively at  $\lambda=1.05 \mu\text{m}$ . However, if required the effect of refractive index variation with the wavelength can also be included. It should be noted that without doping at the C, D, and E layers, the stability increases almost linearly with an increase in the wavelength. The increased wavelength reduces mode confinement and effective index values and this also increases the separation between the modal index values. A similar effect is noticed after the introduction of doping at layers C, D and E layers. When the center wavelength is  $\lambda=1.05 \mu\text{m}$ , the modal stability values are  $S_1=0.0016254$  and  $S_2=0.0017357$ , which are improvements of 35% and 38% from their undoped values, respectively. However, it can be noted that with a change in the wavelength from  $\lambda=1.05 \mu\text{m}$  to  $\lambda=0.85 \mu\text{m}$  the stability difference after doping shows a similar trend to the undoped case. The modal stability values at  $\lambda=0.85 \mu\text{m}$  reduced to  $S_1=0.0010095$  and  $S_2=0.0014572$ , however this reflects a 25% and 74%, improvement from the standard fiber. Similar behavior is observed when the center wavelength is increased to  $\lambda=1.25 \mu\text{m}$  and modal stability increases to  $S_1=0.002199$  and  $S_2=0.0019263$ , and these reflect improvements 35% and 25% increase from the standard fiber. The modal stability still remains 25% higher for the wavelength range from  $\lambda=0.85$

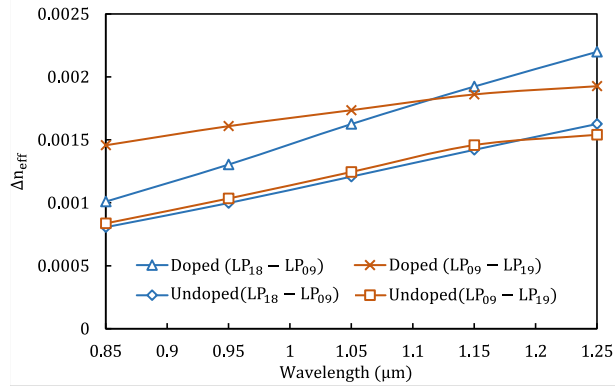


Fig. 8. Effect on the  $\Delta n_{eff}$  with the change in wavelength ( $\lambda$ )

$\mu\text{m}$  to  $\lambda=1.25 \mu\text{m}$ , and hence the proposed design results in a sufficiently improved stability between the competing modes and is also expected to provide for large effective mode areas  $A_{eff}$  because of operating in HOMs [12].

#### 4. Conclusions

We have proposed a novel design approach, which has been validated by rigorous numerical analysis, to improve the stability between the modes can be increased by more than 35%. Increasing the  $\Delta n_{eff}$  between modes results in a more stable and mode-mixing resistant operation and thus allows scalability of power in laser applications. It is also shown here that the design is reasonably stable to possible fabrication tolerance such as the position and width of the doped annular sections.